Group Contribution Lattice Fluid Equation of State for CO₂—Ionic Liquid Systems: An Experimental and Modeling Study

Chengna Dai, Zhigang Lei, Wei Wang, Li Xiao, and Biaohua Chen

State Key Laboratory of Chemical Resource Engineering, Beijing University of Chemical Technology, Beijing 100029, China

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The group contribution lattice fluid equation of state (GCLF EOS) was first extended to predict the thermodynamic properties for carbon dioxide (CO_2)—ionic liquid (IL) systems. The group interaction parameters of CO_2 with IL groups were obtained by means of correlating the exhaustively collected experimental solubility data at high temperatures (above 278.15 K). New group parameters between CO_2 and IL groups were added into the current parameter matrix. It was verified that GCLF EOS with two kinds of mixing rules could be used for predicting the CO_2 solubility in ILs, and volume expansivity of ILs upon the addition of CO_2 , as well as identifying the new structure—property relation. Moreover, it is the first work on the measurement of the solubility of CO_2 in ILs at low temperatures (below 278.15 K), manifesting the applicability of predictive GCLF EOS over a wider temperature range. © 2013 American Institute of Chemical Engineers AIChE J, 59: 4399–4412, 2013

Keywords: GCLF EOS, ILs, CO2 solubility, volume expansivity, structure-property relation, low temperatures

Introduction

Due to their unique advantages such as nonvolatility, high solubility, and "designer solvents", ionic liquids (ILs) have been receiving significant attention in separation processes recently, especially for capturing carbon dioxide (CO₂) from natural gas and postcombustion flue gas.^{1,2} In some cases, the separation requirement for CO₂ yield and purity is not so strict that the adsorption process operating above room temperatures suffices. However, in the purification of syngas for Fischer-Tropsch (FT) synthesis and the production of ammonia (NH₃), the CO₂ content in syngas must be decreased down to several ppm level (usually < 20 ppm) before entering the reaction section.³ For this purpose, the world-famous Rectisol process in which methanol is used as the separating agent for capturing CO₂ operates at temperatures as low as 228 K with high CO₂ solubility. Unfortunately, for capturing CO₂ with ILs, no experimental data on CO₂ solubility below 278 K have been reported by far. Therefore, in this work, we decided to measure the CO2 solubility at low temperatures down to 228 K in three common ILs as the representatives, i.e., $[BMIM]^+[BF_4]^-$, $[HMIM]^+[BF_4]^+$, and [HMIM]⁺[PF₆]⁻. Although there is a great variety of ILs, it is tedious and time-consuming to determine the CO₂ solubility in all kinds of ILs through experiments. Thus, for a better

The predictive models for treating the phase equilibria of the systems with ILs include activity coefficient models, equations of state, molecular simulations, neural networks, and so on, among which activity coefficient models and equations of state are commonly used due to the solid thermodynamic foundation and fast calculation speed. However, the activity coefficient models, including regular solution theory (RST),⁸ universal quasichemical functional-group activity coefficients (UNIFAC),^{4,9-11} and conductor-like screening model for real solvent (COSMO-RS)¹²⁻¹⁵ models, are not functions of volume, and also not dependent on pressure. On the other hand, the equations of state, such as nonrandom lattice-fluid equation of state (NLF EOS), 16,17 the series of statistical associating fluid theory (SAFT) EOS¹⁸⁻²⁸ (e.g., perturbed-chain polar SAFT [PCP-SAFT], soft-SAFT, heterosegmented-SAFT, and SAFT- γ), and contribution equation of state (GC EOS),²⁹ would take into account the dependence of phase volume on pressure for calculating the phase equilibria, which is especially important for estimating volume expansivity. For CO₂-IL systems, Karakatsani et al.30 used tPC-PSAFT EOS to predict the marginal vapor pressure of pure ILs and the solubility of CO_2 , CO, N_2 , and CHF_3 in $[BMIM]^+[PF_6]^-$. The results showed good agreement with experimental data. The soft-SAFT EOS^{25,26} was used to predict the solubility of CO₂, H₂, and Xe in three imidazolium-based IL families with different anions, i.e. $[RMIM]^+[BF_4]^-$, $[RMIM]^+[PF_6]^-$, and $[RMIM]^+[Tf_2N]^-$ (R = ethyl, butyl, hexyl, and octyl), at

and thorough understanding on the separation and thermodynamic behaviors of $\rm CO_2$ –IL systems, reliable and applicable predictive models are indispensable. $^{4-7}$

Additional Supporting Information may be found in the online version of this article.

Correspondence concerning this article should be addressed to Z. Lei at leizhg@mail.buct.edu.cn.

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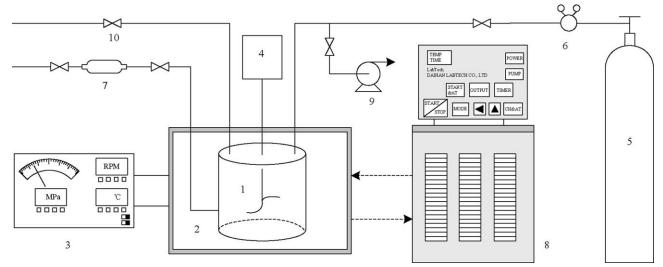


Figure 1. Schematic diagram of the experimental apparatus for measuring the solubility of CO₂ in ILs at low temperatures.

(1) stainless steel autoclave; (2) ethanol bath; (3) pressure and temperature display; (4) stirring paddle; (5) CO₂ cylinder; (6) cylinder regulator (7) liquid sampler; (8) refrigeration compressor; (9) vacuum pump; (10) valve.

pressures up to 100 MPa, and the predicted results were in quantitative agreement with the experimental data. Ji and Adidharma^{22,23} compared the experimental data with the predicted results by heterosegmented-SAFT EOS on the solubility of CO₂ in [RMIM]⁺[BF₄]⁻, [RMIM]⁺[PF₆]⁻, and [RMIM]⁺[Tf₂N]⁻ at temperatures up to 423 K and pressures up to 200 bar, and found that both agreed well, especially at low pressures. The ePC-SAFT model²⁷ introducing a Debye-Hückel Helmholtz energy term improved the prediction of gas solubility in imidazolium-based ILs at temperatures from 293 to 450 K and pressures up to 950 bar. Recently, the SAFT-γ approach²⁴ was developed to predict the phase behavior of CO₂-IL systems by using the optimized group parameters. Although in most cases the SAFT models could give good predictive results, a number of pure component parameters and cross-association binary parameters have to be concerned, and their equation forms are much complicated. Thus, the simple and reliable group-contribution lattice-fluid equation of state (GCLF EOS), which assumes that the interaction energy between particular groups will be constant regardless of the overall structure of the components, is worth developing for predicting the phase equilibrium of CO₂-IL systems.

GCLF EOS was originally developed by High and Danner. 31,32 Later, Lee and Danner extended it to the polymer–solvent systems. 33 A complete parameter matrix of GCLF EOS for both conventional solvents and polymers has already been established. 34 Afterward, Hamedi et al. 35 extended the predictive model to the solubility prediction of CO₂–polymer binary systems. In our previous work, GCLF EOS was firstly extended to systems containing ILs. 36 However, the prediction of gas (especially CO₂) solubility in ILs by GCLF EOS has not been studied so far.

Therefore, the focus of this work is on addressing the cogently interesting issues for GCLF EOS as to (1) adding the group parameters of CO_2 and ILs groups to the current parameter matrix; (2) checking the applicability of GCLF EOS when extrapolated from high (above 278.15 K) to low temperatures (below 278.15 K); (3) checking the availability in predicting the volume expansivity upon addition of CO_2

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into ILs; and (4) identifying the new structure–property relation between molecular structure of ILs and separation performance. For this purpose, the solubility data of CO_2 in $[N_{1,8,8,8}]^+[Tf_2N]^-$, $[N_{1,4,4,4}]^+[Tf_2N]^-$, and $[P_{1,4,4,4}]^+[Tf_2N]^-$ at 298.2 K were measured in this work. The meanings of abbreviations for cations and anions of ILs throughout this article are given in Supporting Information.

Experimental Section

Materials

The ILs $[BMIM]^+[BF_4]^-$, $[HMIM]^+[BF_4]^-$, $[HMIM]^+$ $[PF_6]^-$, $[N_{1,8,8,8}]^+[Tf_2N]^-$, $[N_{1,4,4,4}]^+[Tf_2N]^-$, and $[P_{1,4,4,4}]^+$ $[Tf_2N]^-$ with mass fraction purity > 99 wt % were purchased from Shanghai Chengjie Chemical Co. Before the solubility measurements, the ILs were dried in the vacuum rotary evaporator at 333.2 K for 24 h to remove traces of water and other volatile impurities. The water content in ILs after drying (but before solubility experiments) was less than 400 ppm as determined by Karl Fischer titration (SC-6) in our laboratory, but it became less than 700 ppm after solubility experiments. Thus, the recycled ILs should be dried again to the same water content before further use. A high purity of CO_2 (>99.995 wt %) supplied by Longkou City Gas Plant was used without further purification.

Apparatus and procedure

The solubility of CO_2 in ILs was measured using a low-temperature equilibrium technique, and the apparatus is shown schematically in Figure 1, which mainly consists of a 500 mL stainless steel equilibrium cell, a temperature-controlled ethanol bath producing low temperature down to 213 K in connection with a refrigeration compressor (Model LCC-R220U, made by Daihan Labtech Co.), a stirring paddle, and a liquid sampler.

A certain amount of IL (about 200 mL) was first loaded into the equilibrium cell that was immersed into the ethanol bath for controlling the cooling temperature with a fluctuation of \pm 0.1 K. The temperature was monitored with a thermocouple and a digital thermometer with an uncertainty of 0.1 K. Then

the air in the equilibrium cell was purged out of the system by vacuum pump (2XZ-1). Afterward, a charge of CO_2 was introduced into the equilibrium cell. When CO_2 was dissolved into the IL, the system pressure decreased gradually. The IL phase was mixed with a stirring paddle at a stirring speed of about 200 rpm. It was assumed that gas-liquid equilibrium had been reached until the system pressure measured by a pressure gauge (HQ sensor 1000) with an uncertainty of \pm 0.01 MPa was invariable. Generally, the gas-liquid equilibrium took about three days for each data point.

The gas phase was considered as pure CO_2 gas (i.e., $y_1 = 0$) since IL had a negligible vapor pressure,³⁷ and thus no need to analyze. For each run, a small amount of liquid sample about 2 mL was taken out from the liquid phase. After sampling, the pressure in the equilibrium cell was kept constant with a fluctuation of 0.01 MPa by charging additional CO₂ for next reproducible measurement. Meanwhile, the amounts of CO2 and IL in the liquid sample were determined using gravimetric method by measuring the weight of solutions with and without CO2 in the similar way as reported in previous publication.³⁸ The mass uncertainty of the electronic balance (CPA 1003S, Sartorius) was 0.001 g. Therefore, the estimated uncertainty of solubility measurement in mole fractions was less than 0.006. For each data point, three runs were performed to check the reproducibility of the results at the same temperature and pressure, and to ensure that no degradation of IL took place. The reliability of experimental apparatus and procedure was validated by comparing the obtained experimental solubility data of CO₂ at 298.2 K in [BMIM]⁺[BF₄]⁻ and [HMIM]⁺[BF₄]⁻ with the results reported in different references, ^{16,39–43} and it was found that the ARDs (average relative deviations) were 2.29% for $[BMIM]^+[BF_4]^-$ (Ref. 40) and 1.82% for $[HMIM]^+[BF_4]^-$ (Ref. 16), respectively (see Supporting Information for detailed values).

Thermodynamic Model

Model description

In this work, ILs were divided into several separate functional groups, and the skeleton of cation and anion was treated as a whole functional group as the same manner as in the UNIFAC model and other group-contribution equations of state for ILs. 4,16,36,44,45 Thus, the additional Debye–Hückel term accounting for long-range (LR) electrostatic contributions can be avoided.

In the GCLF EOS, each molecule or group is assumed to occupy a number of lattice sites, and each lattice site is assumed to have a constant volume ($v_h = 9.75 \times 10^{-3} \cdot \text{m}^3 \cdot \text{kmol}^{-1}$) with a fixed coordination number (z = 10) as commonly used. ^{31–36} The equation form of GCLF EOS is written as

$$\frac{\tilde{P}}{\tilde{T}} = \ln\left(\frac{\tilde{v}}{\tilde{v} - 1}\right) + \frac{z}{2}\ln\left(\frac{\tilde{v} + q}{r - 1\tilde{v}}\right) - \frac{\theta^2}{\tilde{T}}$$
(1)

where \tilde{P} , \tilde{T} , and \tilde{v} are the reduced pressure, temperature, and molar volume, respectively, and expressed as

$$\tilde{P} = \frac{P}{P^*} = \frac{2Pv_h}{z\varepsilon^*}, \quad \tilde{T} = \frac{T}{T^*} = \frac{2RT}{z\varepsilon^*}, \quad \tilde{v} = \frac{v}{v^*} = v_h r$$
 (2)

$$\theta = \frac{q/r}{\tilde{v} + q/r - 1}, \quad zq = (z - 2)r + 2 \tag{3}$$

where q is the interaction surface area parameter; r is the number of lattice sites occupied by a molecule or a group;

and R is the universal gas constant (8.314 J·mol⁻¹). The adjustable parameters contain the molecular interaction energy ε^* and molecular reference volume v^* .

For a pure component, GCLF EOS contains only two adjustable parameters, ε^* and v^* , as mentioned above. Once the two parameters are known, all of the remaining parameters in Eq. 1 can be determined from Eqs. 2 and 3 at a given temperature and pressure, and then the properties of a system can be obtained according to Eq. 1 with respect to reduced volume. ε^* and v^* are obtained from the following mixing rule

$$\varepsilon_i^* = \sum_k \sum_m \Theta_k^{(i)} \Theta_m^{(i)} (e_{kk} e_{mm})^{1/2}$$
 (4)

$$v_i^* = \sum_k n_k^{(i)} R_k \tag{5}$$

where e_{kk} is the group interaction energy between like groups k; $\Theta_k^{(i)}$ is the surface area fraction of group k in pure component i; R_k is the group reference volume; and $n_k^{(i)}$ is the number of group k in pure component i

$$e_{kk} = e_{0,k} + e_{1,k} \left(\frac{T}{T_0}\right) + e_{2,k} \left(\frac{T}{T_0}\right)^2$$
 (6)

$$\Theta_k^{(i)} = \frac{n_k^{(i)} Q_k}{\sum_n n_n^{(i)} Q_n}$$
 (7)

$$R_{k} = \frac{1}{10^{3}} \left[R_{0,k} + R_{1,k} \left(\frac{T}{T_{0}} \right) + R_{2,k} \left(\frac{T}{T_{0}} \right)^{2} \right]$$
(8)

where $e_{i,k}$ and $R_{i,k}$ are constants; T (K) is the system temperature; T_0 is arbitrarily set to 273.15 K; and Q_k is the dimensionless surface area parameter of group k, as used in the UNIFAC model.

The densities of most of pure ILs vary from 0.9 to 1.7 g·cm⁻³, which are not very sensitive to temperature change. ⁴⁶ Therefore, for IL groups, it is reasonable to assume that $e_{1,k} = e_{2,k} = 0$ and $R_{1,k} = R_{2,k} = 0$.

For a binary mixture, GCLF EOS also requires one binary interaction parameter k_{12} . The solving procedure is similar to that of pure components, but the following mixing rules are introduced

$$\varepsilon_{i}^{*} = \bar{\theta}_{1} \varepsilon_{11} + \bar{\theta}_{2} \varepsilon_{22} - \bar{\theta}_{1} \bar{\theta}_{2} \dot{\Gamma}_{12} \Delta \varepsilon, \Delta \varepsilon = \varepsilon_{11} + \varepsilon_{22} - 2\varepsilon_{12} \tag{9}$$

$$\varepsilon_{12} = (\varepsilon_{11}\varepsilon_{22})^{1/2}(1 - k_{12}), \varepsilon_{ii} = \sum_{k} \sum_{m} \Theta_{k}^{(i)} \Theta_{m}^{(i)} (e_{kk}e_{mm})^{1/2}$$
 (10)

$$v^* = \sum_i x_i v_i^* \tag{11}$$

where $\dot{\Gamma}_{12}$ is the nonrandomness parameter between molecules 1 and 2. The quasichemical approach gives the following relationship among the nonrandomness parameters

$$\frac{\dot{\Gamma}_{11}\dot{\Gamma}_{22}}{\dot{\Gamma}_{12}^2} = \exp\left(\theta \frac{\Delta \varepsilon}{RT}\right) \tag{12}$$

Table 1. Group Binary Interaction Parameters $(\alpha_{12} = \alpha_{21})$ Between CO₂ (1) and IL (2) Groups for Mixing Rules MR1 and MR2

	α	12
Main Groups	MR1	MR2
[MIM][BF ₄]	-0.3440	-0.0319
[MIM][Tf ₂ N]	-0.5956	-0.1728
[MIM][PF ₆]	-0.3840	-0.0476
[MIM][TfO]	-0.5329	-0.1328
[MIM][MeSO ₄]	-0.2314	-0.0724
[MIM][EtSO ₄]	-0.3261	-0.0877
[MIM][Cl]	-0.1933	0.0348
[MIM][DEPO ₄]	-0.4051	-0.0939
$[MIM][DMPO_4]$	-0.2260	-0.0506
[MIM][MDEGSO ₄]	-0.4063	-0.0961
$[MIM][NO_3]$	-0.4810	-0.0468
[MIM][SCN]	-0.4287	-0.1164
[MIM][TFA]	-0.4349	-0.0881
$[OCH_2MIM][Tf_2N]$	-3.2192	-0.7049
$[MPY][BF_4]$	-0.4105	-0.0080
$[MPY][Tf_2N]$	-0.6032	-0.1677
[MPYR][Tf ₂ N]	-0.5734	-0.1527
[MPYR][TfO]	-0.4762	-0.0849
$[N][Tf_2N]$	-0.9433	-0.1228
[P][Tf ₂ N]	0.2771	-0.3216
[P][Cl]	10.6343	-0.2728
[P][TOS]	1.6997	0.3688

The interaction parameter k_{12} can be calculated by two kinds of mixing rules, namely as MR1 and MR2, expressed

$$k_{12} = \sum_{m} \sum_{n} \Theta_{m}^{(M)} \Theta_{n}^{(M)} \alpha_{mn}$$
 (13)

$$\Theta_k^{(M)} = \frac{\sum_{i} n_k^{(i)} Q_k}{\sum_{p} \sum_{i} n_p^{(i)} Q_p}$$
 (14)

where $\Theta_k^{(M)}$ is the surface area fraction of group k in the mixture containing all groups; and α_{mn} is the group binary interaction parameter. MR1 is expressed in such a way that the interaction parameter is obtained as a correction for the interaction of same and different species. However, for CO₂–polymer systems, Hamedi et al.³⁵ proposed an alternative formulation (MR2) in which the interaction parameter was calculated only as a contribution of groups from different species

$$k_{12} = \sum_{m} \sum_{n} \Theta_{m}^{(1)} \Theta_{n}^{(2)} \alpha_{mn} \text{ (MR2)}$$
 (15)

$$\Theta_k^{(i)} = \frac{n_k^i Q_k}{\sum_p n_p^{(i)} Q_p}$$
 (16)

where $\Theta_k^{(i)}$ is the surface area fraction of group k in pure

Other parameters for a binary mixture are calculated from the following equations

$$r = \sum x_i r_i, \quad q = \sum x_i q_i, \quad \theta = \sum x_i \theta_i$$
 (17)

$$r_i = \frac{v_i^*}{v_i}, \quad zq_i = (z-2)r_i + 2$$
 (18)

$$\theta_i = \frac{q_i/r_i}{\tilde{v}_i/r_i - r_i + q_i}, \quad \bar{\theta}_i = \frac{x_i q_i}{q}$$
 (19)

$$\bar{\theta}_1\dot{\Gamma}_{11} + \bar{\theta}_2\dot{\Gamma}_{12} = \bar{\theta}_2\dot{\Gamma}_{22} + \bar{\theta}_1\dot{\Gamma}_{12} = 1 \tag{20}$$

In GCLF EOS, mole fraction activity coefficient (MFAC) of component i in the mixture is given as follows

$$\ln \gamma_{i} = \ln \frac{a_{i}}{x_{i}} = \ln \varphi_{i} - \ln x_{i} + \ln \tilde{v}_{i} \tilde{v} + q_{i} \ln \left(\frac{\tilde{v}}{\tilde{v} - 1} \frac{\tilde{v}_{i} - 1}{\tilde{v}_{i}} \right) + q_{i} \left(\frac{2\theta_{i,p} - \theta}{\tilde{T}_{i}} - \theta \tilde{T} \right) + \frac{zq_{i}}{2} \ln \dot{\Gamma}_{i}$$
(21)

$$\varphi_{i} = \frac{x_{i}v_{i}^{*}}{\sum_{i} x_{j}v_{j}^{*}} = \frac{x_{i}r_{i}}{\sum_{i} x_{j}r_{j}}$$
(22)

where the subscript i represents a pure component i; x_i is the mole fraction of component i in the liquid phase; Γ_i is nonrandomness parameter; φ_i is the volume fraction of component i in the mixture; and $\theta_{i,p}$ is the surface area fraction of pure component i at the same temperature and pressure as in

The group parameters $(e_{0,k}, R_{0,k}, Q_k)$ for IL groups, and $(e_{0,k}, e_{1,k}, e_{2,k}, R_{0,k}, R_{1,k}, R_{2,k}, Q_k)$ for CO₂ and other groups have been reported in previous references. ^{34–36,47–49} Therefore, the group binary interaction parameters (α_{mn}) between CO₂ and IL groups are the only required input parameters for solving the above equations, which were derived by means of correlating the experimental solubility data.

Procedure of the estimation of group interaction parameters

The following minimized objective function (OF) was used to obtain the group binary interaction parameters α_{mn} between CO₂ and IL groups

$$OF = \sum_{i=1}^{N} \left| \frac{x_{i,IL}(GCLF \quad EOS) - x_{i,IL}(exp.)}{x_{i,IL}(exp.)} \right|$$
(23)

where N is the number of data points; $x_{i,IL}$ (exp.) is the experimental solubility of CO2 in ILs exhaustively collected from references by the end of December 2012; and $x_{i,IL}$ (GCLF EOS) is the predicted result by GCLF EOS. The kinds of ILs, more than 5300 data points, experimental methods, and the corresponding cited references are provided in detail in Supporting Information. It should be noted that all experimental temperatures reported in the references were above 278.15 K. The SOLVER function with the optimization algorithm of Newton's central difference in Microsoft Excel 2003 was used to correlate α_{mn} to minimize the objective function. In this way, we can clearly deal with so many experimental data collected from references. As a result, the obtained group binary interaction parameters α_{mn} between CO2 and 22 IL groups using mixing rules MR1 and MR2 are listed in Table 1. The current GCLF EOS parameter matrix is also illustrated in Figure 2. It should be mentioned that the group binary interaction parameters which are available in previous references remain constant.

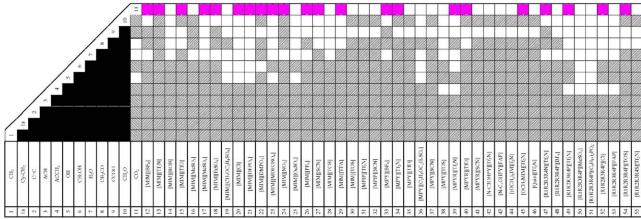


Figure 2. Current GCLF EOS parameter matrix for IL groups.

■ Previously published parameters⁴⁷;
Previously published parameters⁸;
Previously published parameters⁸;
New parameters;
No parameters available. [Color figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

Results and Discussion

Prediction of solubility of CO₂ in ILs at high temperatures

The solubility data of CO_2 in 45 kinds of ILs were investigated at temperatures of 278–453 K and at pressures of 0.01–971.00 bar, and the comparison of experimental solubility of CO_2 in ILs with the predicted results (MR1 and MR2) by GCLF EOS is summarized in Table 2, where the ARD is defined as

ARD (%) =
$$\frac{1}{N} \sum_{i=1}^{N} \left| \frac{x_{i,IL}(\text{cal.}) - x_{i,IL}(\text{exp.})}{x_{i,IL}(\text{exp.})} \right| \times 100$$
 (24)

The predicted CO₂ solubility data in five common ILs, i.e., [BMIM]⁺[BF₄]⁻, [BMIM]⁺[PF₆]⁻, [BMIM]⁺[Tf₂N]⁻, [BMPYR]⁺[Tf₂N]⁻, and [BPY]⁺[BF₄]⁻, are exemplified in Figure 3, along with the experimental results obtained from references. They exhibit the similar trend and agree very well. The extensive experimental data and the predicted results by GCLF EOS with mixing rules MR1 and MR2, as well as those by the COSMO-RS (conductor-like screening model for real solvents) model using the commercial ADF software, ¹⁰⁴ are given in the form of MS Excel spreadsheet in Supporting Information. The comparison shows that GCLF EOS gives better prediction than COSMO-RS model in most cases, and the predicted results by GCLF EOS with both MR1 and MR2 mixing rules are in good agreement with the experimental data.

On the other hand, other equations of state have also been proposed for predicting the solubility of CO_2 in ILs. Ashrafmansouri and Raeissi²⁴ used the SAFT- γ model involving a number of binary interaction parameters (two parameters accounting for the association between CO_2 and anion, two temperature-independent parameters between CO_2 and CH_3 or CH_2 groups, and two temperature-dependent parameters between CO_2 and anion) to predict the solubility of CO_2 in imidazolium-based ILs. The ARDs for $[BMIM]^+[Tf_2N]^-$ and $[HMIM]^+[Tf_2N]^-$ in their work are 2.33% and 4.04%, respectively, while in this work the corresponding ARDs (MR1) are 4.60% (5.26% for MR2) and 5.20% (7.07% for MR2), respectively. Ji et al.²⁷ investigated the density of ILs

and gas solubility in imidazolium-based ILs using ePC-SAFT, in which the IL was considered to be completely dissociated into a cation and an anion, and five parameters for each ion were required. The introduction of Debye–Hückel Helmholtz energy term shows the best performance, with the result that the ARDs of CO₂ solubility in [BMIM] $^+$ [Tf₂N] $^-$ and [EMIM] $^+$ [Tf₂N] $^-$ are 8.10% and 7.60%, respectively, while in this work the corresponding ARDs are 8.50% (8.62% for MR2) and 9.93% (12.28% for MR2), respectively. This demonstrates that the prediction accuracy of GCLF EOS is comparable with other equations of state which require more model parameters. Therefore, GCLF EOS is applicable for predicting the solubility of CO₂ in ILs at high temperatures.

Prediction of solubility of CO₂ in ILs at low temperatures

The solubility data of CO₂ in [BMIM]⁺[BF₄]⁻, [HMIM]⁺ [BF₄]⁻, and [HMIM]⁺[PF₆]⁻ were measured at temperatures (228-273) K and pressures up to 3.0 MPa by means of a low-temperature equilibrium technique. The experimental results are illustrated in Figures 4-6, along with the predicted results by GCLF EOS in which the group binary interaction parameters listed in Table 1 are extrapolated from high to low temperatures. It can be seen that low temperature is favorable for increasing the solubility of CO₂ in ILs significantly. In practice, a high CO₂ solubility at low temperatures can counteract the viscosity increase of pure IL in the mixture. Moreover, the P-T diagrams at constant CO₂ composition in three kinds of ILs ([BMIM]⁺[BF₄]⁻, [HMIM]⁺[BF₄]⁺, and [HMIM]⁺[PF₆]⁻) are plotted in Supporting Information including both low- and hightemperature data, and a "continuation" of CO₂ solubility from high to low temperatures is observed.

In summary, the predicted and experimental results are in good agreement. For the GCLF EOS with mixing rule MR1, the ARDs of CO₂ solubility are 5.98, 6.81, and 12.40% for [BMIM]⁺[BF₄]⁻, [HMIM]⁺[BF₄]⁻, and [HMIM]⁺[PF₆]⁻, respectively, while for the GCLF EOS with mixing rule MR2, they are 6.85, 9.76, and 8.70%, respectively. See Supporting Information for more details. Therefore, the application of GCLF EOS can be extended from high to low temperatures. However, it is worth mentioning that the

Table 2. Comparison of Experimental CO2 Solubility in ILs with the Predicted Results by GCLF EOS

ILs	T range (K)	P range (bar)	MR1 ARD (%)	MR2 ARD (%)	Number of Data Points	Reference
[BMIM] ⁺ [BF ₄] ⁻	303.72–344.49	0.18-0.84	27.70	25.10	21	[50]
Divinivi) [DI 4]	298.20–333.30	12.09-85.00	15.43	16.48	20	[40]
	283.15–323.15	0.02-13.00	14.38	12.25	84	
						[41]
	278.47–368.22	5.87–676.20	17.18	18.08	104	[51]
	282.75-348.15	0.10-20.00	18.00	16.59	36	[39]
	303.93-344.27	0.22-0.92	35.24	32.93	11	[52]
	307.55-322.15	17.50-90.00	10.23	12.74	40	[53]
	298.00-298.20	0.10-20.00	23.84	21.61	9	[42]
	293.25-383.15	10.50-246.00	18.39	21.35	59	[54]
	298.15-298.15	6.50-60.70	10.17	10.17	7	[43]
		total ARD	17.14	17.29	391	
BMIM] ⁺ [Cl] ⁻	353.15-373.15	24.54–369.46	18.36	18.36	45	[55]
BMIM] ⁺ [MDEGSO ₄] ⁻	313.31–333.36	14.30–91.20	32.26	37.45	11	[56]
BMIM] [MeSO ₄]	293.20–413.10	9.08–98.05	11.96	11.83	54	[57]
			23.96	23.96	21	
$BMIM]^{+}[NO_3]^{-}$	313.15–333.15	15.47–93.17				[58]
	298.20-333.20	10.31-88.92	6.78	6.78	17	[40]
	293.13-368.24	3.68-128.32	14.49	14.49	66	[59]
		total ARD	15.14	15.14	104	
$BMIM]^{+}[PF_{6}]^{-}$	313.15-323.15	15.17-95.67	14.09	14.05	21	[58]
	283.15-323.15	0.03-12.99	5.72	5.71	158	[60]
	293.15–393.15	1.05-96.85	17.84	17.88	43	[61]
	298.20–333.40	5.60–146.39	9.85	9.85	70	[40]
	283.15–323.15	0.03-13.00	5.40	5.39	158	
						[41]
	293.29–363.54	4.30–735.00	18.76	18.78	99	[62]
	282.05-348.25	0.10-20.00	12.02	12.01	35	[39]
	297.56–322.52	7.90-80.80	7.34	7.34	42	[63]
	298.15-298.15	5.29-6.67	14.36	14.30	4	[16]
	283.15-343.04	0.41-0.92	36.53	36.50	14	[64]
	298.00-298.20	0.11-20.00	14.84	14.79	9	[42]
	298.15-298.15	2.60-40.20	10.86	10.80	9	[17]
		total ARD	10.42	10.41	662	[]
BMIM] ⁺ [SCN] ⁻	292.35-384.15	10.50–315.00	7.24	7.24	56	[54]
$BMIM]^+[Tf_2N]^-$	298.10–333.30	11.38–132.43	2.68	3.05	55	[40]
DIVITIVI] [1121N]					96	
	283.15–323.15	0.01–13.00	13.53	13.58		[41]
	279.98-339.97	2.92-48.00	7.28	7.57	16	[65]
	283.36–343.78	0.68-1.12	15.45	14.35	14	[66]
	293.35-344.55	10.70-428.00	11.02	9.91	84	[67]
	313.15-453.15	4.20-142.61	7.47	8.78	133	[68]
	292.65-363.26	6.29-499.90	12.75	11.89	68	[69]
	313.20-323.20	80.80-199.40	5.76	5.35	8	[70]
		total ARD	9.73	9.80	474	£ 3
BMIM] ⁺ [TFA] ⁻	298.17-333.41	11.70–92.60	9.69	5.82	19	[56]
Divinivij [1171]	298.00–298.20	0.11–20.00	33.77	20.54	9	[42]
	293.25–363.18	9.79–629.89	21.75	19.35	52	
	293.23-303.16					[71]
D) 40 41±17701=	200 20 222 20	total ARD	20.24	16.27	80	F403
BMIM] ⁺ [TfO] ⁻	298.20-333.30	10.44-114.77	10.42	13.46	27	[40]
	303.85-344.55	8.50-375.00	19.84	16.69	65	[72]
	303.20-343.20	2.15-65.21	20.50	27.14	35	[73]
		total ARD	18.02	18.88	127	
$[MPYR]^{+}[Tf_{2}N]^{-}$	283.15-323.15	0.02-13.00	23.27	27.74	52	[41]
3 2 2 3	303.78-344.15	0.49-0.57	47.22	44.85	11	[74]
	293.10–413.20	2.80–108.13	4.30	4.66	26	[75]
	303.15–373.15	6.80–627.70	12.61	12.09	72	[76]
	313.20–323.20	80.60–200.60	9.18	9.10	8	[70]
	313.20-323.20		16.70			[/0]
DMDVD1+(ECO1-	202.15.272.25	total ARD		17.75	169	1003
BMPYR] ⁺ [TfO] ⁻	303.15–373.25	18.80-702.00	19.18	19.18	64	[77]
BPY] ⁺ [BF ₄] ⁻	313.15–333.15	15.47-95.80	3.24	3.24	21	[58]
$C_2OMIM]^+[Tf_2N]^-$	303.15-323.15	0.10-1.60	23.58	23.58	18	[78]
$C_5 \text{MIM}$] ⁺ $[\text{Tf}_2 \text{N}]$ ⁻	293.30-363.29	6.18-598.05	10.00	10.00	144	[79]
$DMIM]^{+}[Tf_2N]^{-}$	298.15-343.15	14.74-145.96	4.52	11.31	22	[80]
1 1 2 3	298.15-343.15	14.74-145.96	4.52	11.31	22	[81]
	313.15–313.15	20.00–144.00	15.68	21.17	14	[82]
	313.20–323.20	80.70–201.50	9.81	14.59	8	[70]
materia	200 15 200 15	total ARD	7.53	13.80	66	
$\mathrm{EMIM]}^{+}\mathrm{[BF_4]}^{-}$	298.15–298.15	2.51-8.75	26.88	5.19	9	[16]
	303.20-343.20	4.97-43.29	88.77	65.00	25	[83]
	298.20-313.20	5.30-40.60	50.50	29.79	17	[38]
		total ARD	65.09	42.71	51	
EMIM] ⁺ [DEPO ₄] ⁻	313.15-333.15	0.24-1.99	4.44	4.44	22	[84]
[EMIM] EtSO ₄]	313.15–333.15	14.36–94.61	13.92	13.92	21	[58]
	303.20–343.20	2.11–49.62	14.28	14.28	35	
						[73]
	303.15–353.15	1.22–15.47	3.90	3.90	39	[85]
		total ARD	9.94	9.94	95	
EMIM] ⁺ [MDEGSO ₄] ⁻	303.20-343.20	8.54-67.10	5.38	5.38	30	[86]
	308.14-366.03	14.90-971.00	41.54	36.88	74	[87]
$EMIM]^{+}[PF_{6}]^{-}$					8	[16]
	298.15-298 15	2.13-9.03	6.43	14.90	n	
	298.15–298.15	2.13-9.03	6.43	14.90 46.22		
$\mathrm{EMIM}]^+[\mathrm{PF}_6]^- \ \mathrm{EMIM}]^+[\mathrm{Tf}_2\mathrm{N}]^-$	298.15–298.15 303.63–344.23 283.43–343.07	2.13–9.03 0.43–0.57 0.65–0.86	6.43 49.88 22.60	46.22 16.40	14 5	[88] [66]

Table 2. Continued

	T range	P range	MR1 ARD	MR2 ARD	Number of	
ILs	(K)	(bar)	(%)	(%)	Data Points	Reference
	312.10-453.15	6.26–147.70	10.84	15.16	191	[89]
	292.75-344.55	12.20-432.00	10.26	12.02	78	[67]
	297.90-298.20	0.50-20.00	5.19	7.62	9	[42]
	292.16-363.55	6.20-478.50	8.65	8.84	153	[79]
	298.15-343.15	12.35-147.94	2.47	3.20	21	[80]
	298.15-343.15	12.35-147.94	2.47	3.20	21	[81]
		total ARD	11.49	13.39	514	
[EMIM] ⁺ [TFA] ⁻	298.10-298.10	0.10-20.00	10.36	10.15	9	[42]
	298.10-348.20	0.10-20.00	19.89	19.29	27	[90]
		total ARD	17.51	17.01	36	
[EMIM] ⁺ [TfO] ⁻	303.85-344.55	8.00-378.00	12.46	12.54	55	[72]
	303.20-343.20	1.80-58.84	30.18	27.00	30	[91]
		total ARD	18.71	17.64	85	
$[HMIM]^+[BF_4]^-$	293.18-368.16	5.40-866.00	20.56	18.99	104	[92]
3 1 43	298.15-298.15	3.12-8.99	36.67	30.01	8	[16]
	307.55-322.15	21.30-86.40	7.05	4.55	53	[53]
		total ARD	17.00	14.89	165	£3
$[HMIM]^+[PF_6]^-$	298.31-363.58	6.40-946.00	20.12	19.79	98	[93]
[110]	298.15–298.15	2.96–9.27	20.72	14.93	14	[16]
	2,0.10 2,0.10	total ARD	20.20	19.18	112	[10]
$HMIM]^+[Tf_2N]^-$	298.10-333.30	13.15–115.58	3.96	4.73	28	[40]
111/11/11] [112/11]	298.15–298.15	1.64-8.59	17.15	23.36	9	[16]
	293.15–413.20	6.01–99.11	4.14	2.81	25	[94]
	283.16–323.17	0.01-33.11	13.11	14.10	57	[56]
	288.48–343.20	0.29-0.94	42.71	46.92	11	[95]
	281.90–348.60	0.29-0.94	12.97	15.94	72	
					10	[96]
	298.15–298.15	1.57-8.40	4.55	7.12	90	[97]
	303.85–344.55	14.00–390.00	10.55	13.10		[67]
	297.30–297.40	0.09–19.75	16.18	22.09	10	[42]
	298.15–343.15	8.00–247.08	8.94	9.64	36	[80]
	298.15–343.15	17.93–247.08	9.21	9.86	25	[81]
	278.12–368.44	4.22-143.37	5.92	6.90	123	[98]
		total ARD	9.99	11.68	496	
HMIM] ⁺ [TfO] ⁻	313.23–313.39	14.94-84.23	19.42	12.50	6	[99]
	303.85-344.55	12.50-363.00	20.47	15.46	70	[72]
		total ARD	20.39	15.23	76	
$HMMIM]^{+}[Tf_2N]^{-}$	298.20-333.30	14.97-118.04	10.46	5.67	29	[40]
$[HMPY]^+[Tf_2N]^-$	283.18-323.15	0.01-13.00	17.88	18.16	56	[56]
$HMPYR]^{+}[Tf_{2}N]^{-}$	303.15-373.15	10.60-475.50	10.13	10.21	64	[100]
$MMIM]^{+}[DMPO_{4}]^{-}$	313.15-333.15	0.49-1.75	6.83	6.83	12	[84]
$N_{1,8,8,8}]^+[Tf_2N]^-$	313.20-323.20	80.80-205.60	6.87	13.16	8	[70]
$N_{2,1,1,3}]^+[Tf_2N]^-$	313.22-313.25	11.34-94.66	2.69	3.19	8	[99]
$N_{4,1,1,1}]^+[Tf_2N]^-$	333.23-333.23	15.60-80.90	7.26	5.77	6	[56]
	282.93-343.07	0.36-0.89	44.66	46.37	12	[66]
	313.20-323.20	85.80-196.30	11.36	12.44	8	[70]
		total ARD	17.06	13.81	153	
$N_{4,4,4,1}]^{+}[Tf_2N]^{-}$	298.15-298.15	0.05-5.50	19.29	15.18	11	[41]
NMIM] ⁺ [PF ₆] ⁻	293.15–298.15	8.60–35.40	25.83	21.12	11	[17]
$OMIM]^+[BF_4]^-$	313.15–333.15	15.61–93.73	12.87	8.16	21	[58]
J L 73	307.79–363.29	5.71-858.00	20.33	17.73	100	[101]
	307.55–322.15	41.70–87.20	9.59	5.28	32	[53]
	55,.55 522.15	total ARD	17.06	13.81	153	[55]
$OMIM]^{+}[PF_{6}]^{-}$	313.15-333.15	16.00–92.88	18.00	14.58	21	[58]
$OMIMJ^+[Tf_2N]^-$	298.20–333.30	13.26–114.69	3.16	7.60	22	[40]
	297.55–344.55	6.80–348.00	17.78	21.65	97	[67]
	303.15–353.15	1.12–20.63	27.90	33.32	43	[102]
	505.15-555.15	total ARD	18.48	22.84	162	[102]
OMIM] ⁺ [TfO] ⁻	303.85-344.55	6.80–340.00	20.35	16.06	65	[67]
	303.15–373.15				72	
OMPYR] ⁺ [Tf ₂ N] ⁻		5.10–359.20	15.30	17.46		[100]
P _{1,4,4,4}] ⁺ [TOS] ⁻	323.15–323.15	0.0034-13.00	6.49	5.25	33	[41]
$P_{6,6,6,14}]^+[Cl]^-$	302.55–363.68	1.68-245.70	13.97	13.97	69	[103]
	313.20–323.20	82.10–207.10	6.17	6.17	8	[70]
1+1000	202.00.252.22	total ARD	13.16	13.16	77	F + O # -
$[P_{6,6,6,14}]^+[Tf_2N]^-$	292.88–363.53	1.06–721.85	11.64	11.64	120	[103]
	293.35–375.35	5.30-222.00	9.78	9.78	91	[77]
	313.20-323.20	80.90-201.70	3.90	3.90	8	[70]
		total ARD	15.18	15.18	219	

predicted results by GCLF EOS sometimes do not match the trend of experimental data at low temperatures and high pressures because in this case the equilibrium state is in the vicinity of critical point of CO₂.

Prediction of volume expansivity of ILs

The addition of CO₂ into ILs as well as conventional organic solvents will lead to an increase of the liquid

volume. Volume expansivity is an important physical quantity in determining the liquid volume flow rate along the absorption column so as to avoid liquid flooding occurring in the tray or packing. Besides, volume expansivity with high accuracy is also needed in the solubility measurement using a commercial gravimetric microbalance in which buoyancy effect has to be considered. Otherwise, erroneous or even negative solubility data may arise. ^{60,105} Badilla and

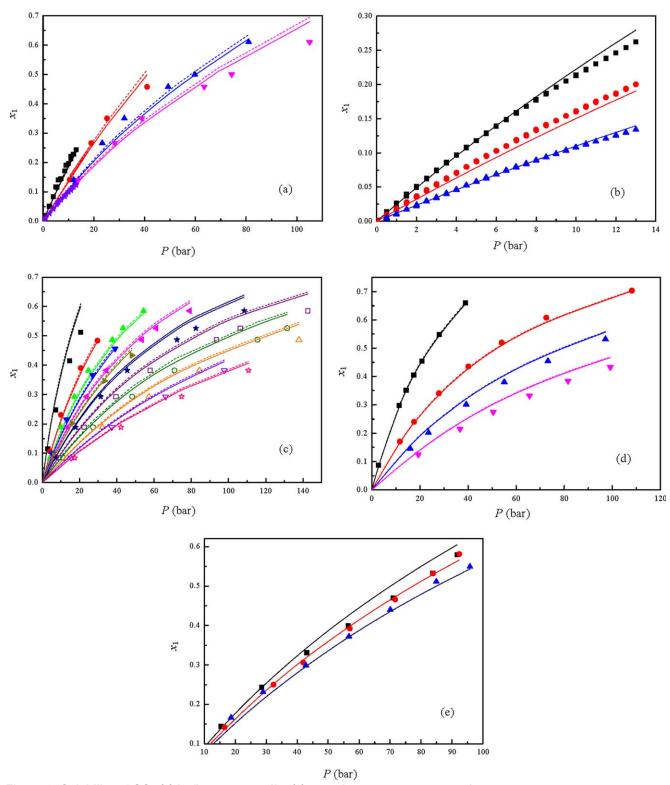


Figure 3. Solubility of CO₂ (1) in five common ILs (2) at various temperatures and pressures.

Solid lines, predicted results by GCLF EOS with MR1; Dashed lines, predicted results by GCLF EOS with MR2; Scattered points, experimental data from references. (a) $[BMIM]^+[BF_4]^-$: () 283.15 K, Ref. 41; () 293 K, Ref. 54; () 303 K, Ref. 54; () 313 K, Ref. 54; () 323.15 K, Ref. 41; () 283.15 K, Ref. 41; () 298.15 K, Ref. 41; () 323.15 K, Ref. 68; () 333.15 K, Ref. 58; () 333

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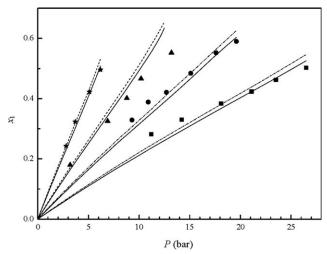


Figure 4. Solubility of CO₂ (1) in [BMIM]⁺[BF₄]⁻ (2) at low temperatures.

Solid lines, predicted results by GCLF EOS with MR1; Dashed lines, predicted results by GCLF EOS with MR2; Scattered points, experimental data. (■) 273.2 K; (●) 258.2 K; (▲) 243.2 K; (★) 228.2 K.

coworkers 106 defined the volume expansivity $\left(\frac{\Delta V}{V}\right)$ of ILs on the basis of molar volume

$$\frac{\Delta V}{V}\% = \frac{\tilde{V}_{\rm M}(T, P, x) - \tilde{V}_{\rm IL}(T, P_0)}{\tilde{V}_{\rm IL}(T, P_0)} \times 100 \tag{25}$$

where $\tilde{V}_{\rm M}(T,P,x)$ is the molar volume of IL/CO₂ mixture at a given temperature (T) and pressure (P); and $\tilde{V}_{\rm IL}(T,P_0)$ is the molar volume of pure IL at the same temperature (T) and ambient pressure ($P_0=1$ bar).

The comparison of volume expansivity for 10 kinds of ILs between the predicted results by GCLF EOS (MR1 and MR2) and experimental data collected from references is summarized in Table 3. The molar volumes of pure ILs and IL/ $\rm CO_2$ mixtures, and the corresponding volume expansivity

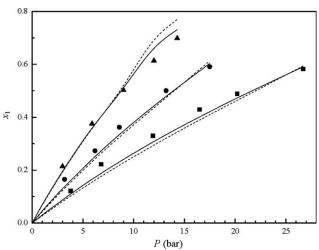


Figure 5. Solubility of CO₂ (1) in [HMIM]⁺[BF₄]⁻ (2) at low temperatures.

Solid lines, predicted results by GCLF EOS with MR1; Dashed lines, predicted results by GCLF EOS with MR2; Scattered points, experimental data. (■) 273.2 K; (●) 258.2 K; (▲) 243.2 K.

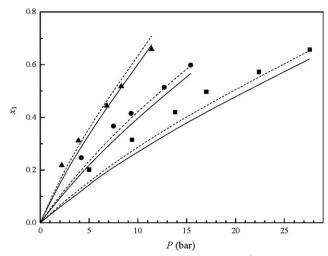


Figure 6. Solubility of CO₂ (1) in [HMIM]⁺[PF₆]⁻ (2) at low temperatures.

Solid lines, predicted results by GCLF EOS with MR1; Dashed lines, predicted results by GCLF EOS with MR2; Scattered points, experimental data. (■) 273.2 K; (●) 258.2 K; (▲) 243.2 K.

are provided in detail in Supporting Information. It can be seen that the predicted results are in good agreement with experimental data, with all the ARDs less than 10%, manifesting the applicability of GCLF EOS for predicting volume expansivity of CO₂–IL systems.

For the imidazolium-based ILs, Aki et al. 40 concluded that volume expansivity was independent on the choice of anions. Herein, we go a further step to explore the influence of different combinations of cations and anions involving 24 kinds of ILs (but not limited to imidazolium-based ILs) and 1261 data points on volume expansivity by using GCLF EOS at temperatures (228–343) K and pressures up to 250 bars. As shown in Figure 7, volume expansivity is almost independent on the kinds of ILs although their molecular weights are quite different. Moreover, there is a linear relationship between the percent volume expansivity $\frac{\Delta V}{V}\%$ and CO₂ solubility in all of the ILs (x)

$$\frac{\Delta V}{V}\% = -89.1950x\tag{26}$$

with the correlation factor $R^2 = 0.9911$. The larger proportionality constant (-89.195) in comparison with conventional

Table 3. Comparison of Volume Expansivity of ILs upon Addition of CO₂ Between Experimental Data and Predicted Results by GCLF EOS

ILs	MR1 ARD (%)	MR2 ARD (%)	Number of Data Points	References
$[BMIM]^+[BF_4]^-$	5.59	5.56	20	[40]
[BMIM] ⁺ [TfO] ⁻	3.85	3.84	27	[40]
$[BMIM]^+[NO_3]^-$	5.03	5.08	17	[40]
$[BMIM]^+[PF_6]^-$	5.12	5.40	70	[40]
$[BMIM]^+[Tf_2N]^-$	1.60	1.60	59	[40]
$[DMIM]^{+}[Tf_2N]^{-}$	4.01	4.00	22	[81]
$[EMIM]^+[Tf_2N]^-$	4.41	4.42	21	[81]
$[HMMIM]^+[Tf_2N]^-$	2.05	1.84	29	[40]
$[HMIM]^+[Tf_2N]^-$	1.31	1.31	28	[40]
$[HMIM]^+[Tf_2N]^-$	9.27	9.26	26	[81]
$[OMIM]^+[Tf_2N]^-$	0.90	0.91	22	[40]

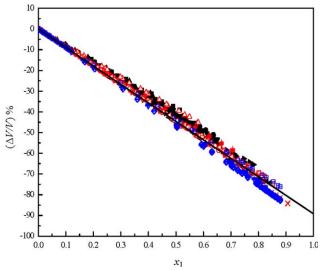


Figure 7. Volume expansivity upon addition of CO₂ (1) for various ILs (2) predicted by GCLF EOS with MR1.

(\blacksquare) [BMIM][BF₄]; (\bullet) [BMIM][TfO]; (\blacktriangle) [BMIM] $[NO_3]; (\blacktriangledown) [BMIM][PF_6]; (\spadesuit) [BMIM][Tf_2N];$ $(\blacktriangleleft) \quad [\mathrm{DMIM}][\mathrm{Tf}_2\mathrm{N}]; \quad (\blacktriangleright) \quad [\mathrm{EMIM}][\mathrm{Tf}_2\mathrm{N}]; \quad (\bigstar) \\ [\mathrm{HMMIM}][\mathrm{Tf}_2\mathrm{N}]; \quad (\Box) \quad [\mathrm{HMIM}][\mathrm{Tf}_2\mathrm{N}]; \quad (\bigcirc)$ $[OMIM][Tf_2N]; (\triangle) [BPY][BF_4];$ $[Tf_2N]; (\lozenge) [HMPY][Tf_2N]; (\triangleleft) [BMPYR][Tf_2N];$ $\begin{array}{l} [T_1_2N]; \; (>) \; [HMPTR][T_1_2N]; \; (>) \; [BMPTR][T_1_2N]; \; (>) \; [MMPYR][Tf_2N]; \; (+) \; [MMPYR][Tf_2N]; \; (+) \; [BMPYR][Tf_2N]; \; (>) \; [N_{1,8,8,8}][Tf_2N]; \; (=) \; [N_{4,1,1,1}] \; [Tf_2N]; \; (>) \; [N_{2,1,1,3}][Tf_2N]; \; (>) \; [N_{1,4,4,4}][Tf_2N]; \; (=) \; [P_{6,6,6,14}][CTf_2N]; \; (=) \; [P_{1,4,4,4}] \; [TOS]; \; Solid lines, linear fitting results. [Color forms on the ries and in the online is an expectation of the control of the control$ figure can be viewed in the online issue, which is available at wileyonlinelibrary.com.]

organic solvents represents a small volume expansion upon addition of CO₂, indicating that the lattice volume and shape do not change substantially before and after the dissolution of CO₂ in ILs due to the strong electrostatic interaction between

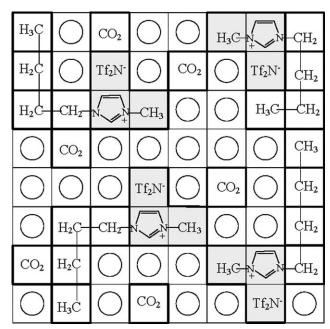


Figure 8. Schematic representation of lattice model for CO₂ molecules, cations, anions, and holes.

cations and anions. Thus, one unique characteristic of ILs is that there is very little volume expansivity upon CO₂ dissolution in the ILs. The solubility difference among ILs will be in some degree attributed to the different fractions of hole volumes (i.e., free volumes) that can be occupied by CO₂ according to the lattice-fluid theory, as illustrated in Figure 8. Therefore, the ILs with long alkyl chain length and weak cation-anion interaction are favorable for increasing CO2 solubility due to the increased free volume. To the best of our knowledge, this is the first work on predicting and analyzing the volume expansivity by using a predictive thermodynamic

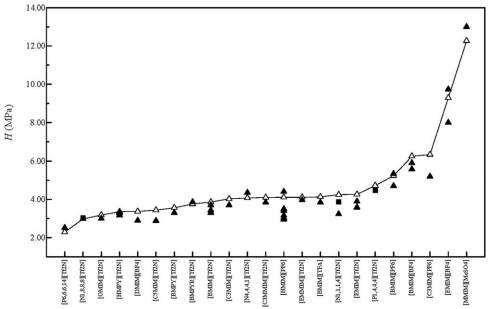


Figure 9. Henry's constants of CO_2 in various ILs at T = 298.15 K.

 (\triangle) predicted results by GCLF EOS with MR1; (\blacktriangle) experimental data from references; $^{38,41,56,64,66,79,95,107-112}$ (\blacksquare) experimental data obtained in this work.

Structure-property relation for the solubility of CO_2 in ILs

One of the most important application of predictive models is to identify the structure–property relation between molecular structure of ILs and separation performance (i.e., Henry's constant which reflects the magnitude of CO₂ solubility in various ILs at low pressures) for physical absorption. Based on the mole fraction of CO₂ in CO₂–IL mixture, the Henry's constant can be written as

$$H(T,P) = \lim_{x \to 0} \frac{yP\phi(T,P)}{x} = \lim_{x \to 0} \gamma P^{0}(T) = \gamma^{\infty} P^{0}(T)$$
 (27)

where x is mole fraction of CO_2 in the liquid phase (i.e., CO_2 solubility); T and P are the system temperature and pressure, respectively; P^0 is the saturated vapor pressure of CO_2 gas at temperature T, obtained from the Antoine equation using the Antoine constants as proposed by Shiflett and Yokozeki¹⁰⁵; and γ^∞ is the activity coefficient of CO_2 in ILs at infinite dilution, calculated by GCLF EOS with MR1 or MR2 mixing rule. Figure 9 shows the Henry's constants of CO_2 in 23 ILs at 298.15 K. It can be seen that the predicted Henry's constants by GCLF EOS (MR1) exhibit the similar trend as the experimental values collected from references. $^{38,41,56,64,66,79,95,107-112}$

In order to achieve a deep insight into structure-property relation, we measured the solubility of CO_2 in $[N_{1.8.8.8}]^+$ $[Tf_2N]^-$, $[N_{1,4,4,4}]^+[Tf_2N]^-$, and $[P_{1,4,4,4}]^+[Tf_2N]^-$ at 298.2 K and pressures up to 8.0 bar in this work, and the solubility data are listed in Supporting Information. The corresponding Henry's constants are deduced from the linear extrapolation of experimental CO₂ solubility in ILs. As shown in Figure 9, for the ILs with the same anion [Tf₂N], the difference of Henry's constants among imidazolium, pyrrolidinium, and pyridinium-based cations is not significant with n (the number of carbon atoms in the alkyl chain on cation of ILs) < 8. However, for the ILs with ammonium and phosphoniumbased cations and the same anion [Tf2N], the Henry's con- $\begin{array}{l} \text{stants} \quad \text{follow} \quad \text{the} \quad \text{order} \quad \text{of} \quad {[P_{6,6,6,14}]}^+ < {[N_{1,8,8,8}]}^+ \\ < {[N_{4,4,4,1}]}^+ < {[N_{1,1,1,4}]}^+ < {[P_{1,4,4,4}]}^+, \quad \text{and} \quad \text{the} \quad \text{maximum} \end{array}$ value with n < 8 is twice as much as the minimum value with n >> 8. Similarly, for the ILs with imidazolium-based cations and the same anion [BF4]-, the maximum Henry's constant for [EMIM]⁺[BF₄]⁻ is triplicate as much as the minimum one for [DMIM]⁺[BF₄]⁻. Moreover, for most of the investigated ILs with the anion [Tf₂N]⁻, their Henry's constants are indeed higher than that for $[DMIM]^+[BF_4]^-$. Therefore, the long alkyl chain length on the cation may influence CO2 solubility effectively. Care should be taken to consider the cation effect, although some authors claimed that the primary contribution to gas solubility is anion effect. 53,113,114

Conclusions

The application of GCLF EOS with two mixing rules, i.e., MR1 and MR2, has been extended to the important CO_2 –IL systems, and 22 new group binary interaction parameters α_{mn} between CO_2 and IL groups were added into the current GCLF EOS parameter matrix so that it may become a universal and available equation of state for conventional organic solvents, polymers, ILs, and gas molecules. From this work, the following conclusions can be drawn: (1)

GCLF EOS can be used for predicting the CO₂ solubility in ILs at either high (above 278.15 K) or low temperatures (below 278.15 K), although the group binary interaction parameters were derived from the experimental solubility data at high temperatures. To the best of our knowledge, this should be the first work that presents the solubility data of CO₂ in ILs at low temperatures (below 278.15 K). (2) The volume expansivity is independent on the different combinations of cations and anions, and exhibits a linear relationship with CO₂ solubility as defined in Eq. 26 over a wide temperature and pressure range. (3) The influence of alkyl chain length on the cation on CO2 solubility cannot be underestimated in some cases. With the same anion, there is a minor difference in Henry's constants for ILs with a short alkyl chain length on the cation. But when the alkyl chain length increases to a certain degree, the Henry's constant will decrease significantly because the free volume effect becomes effective. Anyway, this work confirms the strong predictive power of GCLF EOS for CO2-IL systems.

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